

Helium fine structure theory for the determination of α

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Introduction

- comparison of high precision theoretical predictions of atomic energy levels to experiments gives information on the structure of nuclei: rms charge radius, nuclear polarizability, magnetic moment, etc.
- determining fundamental constants from the atomic and molecular structure
- accurate treatment of electron correlations
- beyond static nucleus: finite nuclear mass corrections including relativistic effects

α from the fine structure in hydrogen

With the inclusion of the finite mass of the nucleus, relativistic, and QED corrections up to order $\alpha^3\nu_0$, the value of the fine structure constant is determined to be

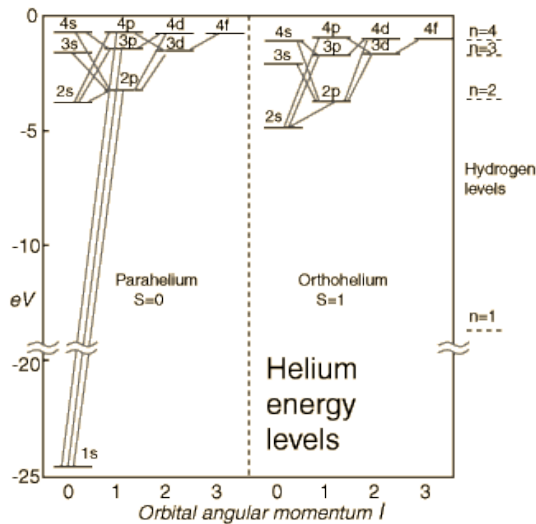
$$\alpha^{-1}(\text{H}) = 137.035\,45(62).$$

It is consistent with the most accurately known value of α at present [Hanneke 2008, Kinoshita 2007]

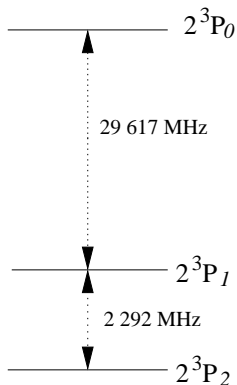
$$\alpha^{-1}(g-2) = 137.035\,999\,084(51).$$

The 4.5 ppm determination of α from hydrogen is much less precise than the 0.37 ppb value from $g-2$ because of the short lifetime of hydrogenic $2p$ states, of order 10^{-9} seconds.

Helium energy levels



Helium fine structure



- the large interval is used for determining α and the small one as a test of theory.

α from the helium fine structure

- determining α from the simplest many-electron atom, helium, program initiated in 1964 by Schwartz.
- the high experimental precision achieved in helium make possible accurate determination of α that depends on the low energy scales characteristic of atomic physics.
- three states is 2^3P_J , nonrelativistically degenerate, but relativistic effects lead to a frequency splitting of order $\alpha^2 R_\infty c$.

Expansion of binding energy in α

Expansion of the energy in powers of the fine structure constant
($\alpha \approx 1/137$)

$$E_{\text{fs}}(\alpha) = E_{\text{fs}}^{(4)} + E_{\text{fs}}^{(6)} + E_{\text{fs}}^{(7)} + \dots$$

- $E^{(n)} \sim \alpha^n \mathcal{E}^{(n)}$
- Valid for small systems with not too large nuclear charge Z
- Expansion coefficients are expressed in terms of nonrelativistic expectation values of effective Hamiltonians
- $E_{\text{fs}}^{(4)} = \langle H_{\text{fs}} \rangle$

Leading relativistic, QED and the nuclear recoil

$$\begin{aligned}
 H_{fs} &= \frac{\alpha}{4m^2} \left(\frac{\vec{\sigma}_1 \cdot \vec{\sigma}_2}{r^3} - 3 \frac{\vec{\sigma}_1 \cdot \vec{r} \vec{\sigma}_2 \cdot \vec{r}}{r^5} \right) (1 + a_e)^2 \\
 &+ \frac{Z\alpha}{4m^2} \left[\frac{1}{r_1^3} \vec{r}_1 \times \vec{p}_1 \cdot \vec{\sigma}_1 + \frac{1}{r_2^3} \vec{r}_2 \times \vec{p}_2 \cdot \vec{\sigma}_2 \right] (1 + 2a_e) \\
 &+ \frac{\alpha}{4m^2 r^3} \left[[(1 + 2a_e) \vec{\sigma}_2 + 2(1 + a_e) \vec{\sigma}_1] \cdot \vec{r} \times \vec{p}_2 \right. \\
 &\quad \left. - [(1 + 2a_e) \vec{\sigma}_1 + 2(1 + a_e) \vec{\sigma}_2] \cdot \vec{r} \times \vec{p}_1 \right] \\
 &+ \frac{Z\alpha}{2mM} \left[\frac{\vec{r}_1}{r_1^3} \times (\vec{p}_1 + \vec{p}_2) \cdot \vec{\sigma}_1 + \frac{\vec{r}_2}{r_2^3} \times (\vec{p}_1 + \vec{p}_2) \cdot \vec{\sigma}_2 \right] (1 + a_e)
 \end{aligned}$$

Higher order corrections

$$E^{(6)} = \langle H^{(6)} \rangle + \langle H^{(4)} \frac{1}{(E_0 - H_0)'} H^{(4)} \rangle$$

$$E^{(7)} = \langle H^{(7)} \rangle + 2 \left\langle H^{(5)} \frac{1}{(E_0 - H_0)'} H_{fs}^{(4)} \right\rangle + E_L$$

$$H^{(5)} = -\frac{7}{6\pi} \frac{\alpha^2}{r^3} + \frac{38Z\alpha^2}{45} [\delta^3(r_1) + \delta^3(r_2)]$$

- anomalous magnetic moment
- electron self-energy and vacuum-polarization
- finite nuclear mass effects

Higher order effective Hamiltonian $H^{(7)}$

$$\begin{aligned} H^{(7)} = & Z \alpha^7 \left(\frac{91}{180} + \frac{2}{3} \ln[(Z \alpha)^{-2}] \right) i \vec{p}_1 \times \delta^3(r_1) \vec{p}_1 \cdot \vec{\sigma}_1 \\ & + \alpha^7 \left(-\frac{83}{60} + \frac{\ln \alpha}{2} \right) (\vec{\sigma}_1 \cdot \vec{\nabla}) (\vec{\sigma}_2 \cdot \vec{\nabla}) \delta^3(r) \\ & - \alpha^7 \frac{15}{8\pi} \frac{1}{r^7} (\vec{\sigma}_1 \cdot \vec{r}) (\vec{\sigma}_2 \cdot \vec{r}) \\ & + \alpha^7 \left(\frac{69}{10} + 3 \ln \alpha \right) i \vec{p}_1 \times \delta^3(r) \vec{p}_1 \cdot \vec{\sigma}_1 \\ & - \alpha^7 \frac{3}{4\pi} i \vec{p}_1 \times \frac{1}{r^3} \vec{p}_1 \cdot \vec{\sigma}_1 + H_{\text{amm}} \end{aligned}$$

- dimensional regularization

Nonrelativistic wave function

- $\vec{\phi}(\vec{r}_1, \vec{r}_2) = \sum_{i=1}^N c_i [\vec{r}_1 \exp(-\alpha_i r_1 - \beta_i r_2 - \gamma_i r) - (1 \leftrightarrow 2)]$
- variational approach: minimize energy with respect to $c_i, \alpha_i, \beta_i, \gamma_i$
- master integral

$$\frac{1}{16\pi^2} \int d^3r_1 \int d^3r_2 \frac{e^{-\alpha r_1 - \beta r_2 - \gamma r_{12}}}{r_1 r_2 r_{12}} = \frac{1}{(\alpha + \beta)(\beta + \gamma)(\gamma + \alpha)}$$

- parameters $\alpha_i, \beta_i,$ and γ_i are chosen quasirandomly

$$\alpha_i \in [A_1, A_2]$$

$$\beta_i \in [B_1, B_2]$$

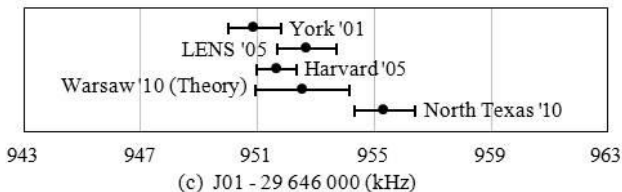
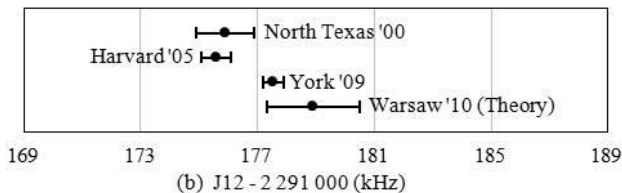
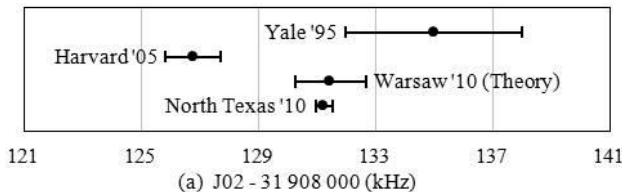
$$\gamma_i \in [C_1, C_2]$$

- $E_0(2^3P) = -2.133\,164\,190\,779\,283\,205\,146\,96_{-10}^{+0}$

Results

Term	ν_{01}	ν_{12}	Ref.
$m\alpha^4(+m/M)$	29 563 765.45	2 320 241.43	Drake (2002)
	29 563 765.23	2 320 241.42	
$m\alpha^5(+m/M)$	54 704.04	-22 544.00	Drake (2002)
	54 704.04	-22 545.01	
$m\alpha^6$	-1 607.52(2)	-6 506.43	Drake (2002)
	-1 607.61(4)	-6 506.45(7)	
$m\alpha^6 m/M$	-9.96	9.15	Drake (2002)
	-10.37(5)	9.80(11)	
$m\alpha^7 \log(Z\alpha)$	81.43	-5.87	Drake (2002)
	81.42	-5.87	
$m\alpha^7, \text{nlog}$	18.86	-14.38	
$m\alpha^8$	± 1.7	± 1.7	
Total theory	$29\,616\,952.29 \pm 1.7$	$2\,291\,178.91 \pm 1.7$	

Comparison with experiments: from Shiner 2010



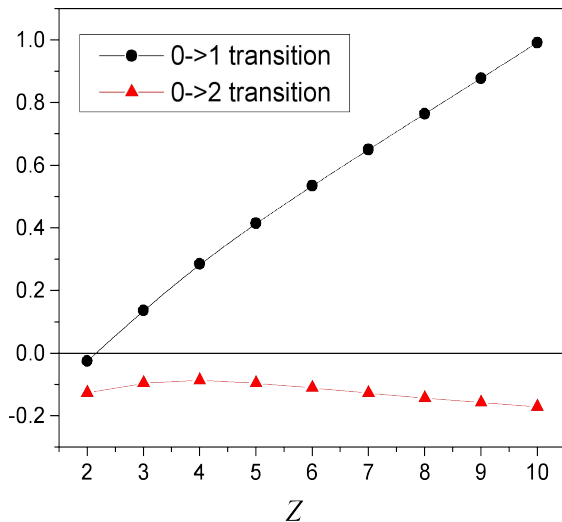
Determination of the fine structure constant

- α from ν_{01} Zelevinsky, Farkas, and Gabrielse (2005)

$$\alpha^{-1}(\text{He}) = 137.036\,001\,1(39)_{\text{theo}}(16)_{\text{exp}}$$

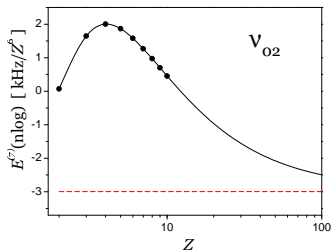
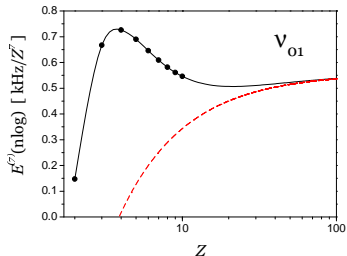
- $\alpha^{-1}(g-2) = 137.035\,999\,084(51)$
- theoretical uncertainty is due to the higher order terms
- ν_{02} fine structure measurement in heliumlike ions ?

$m\alpha^6$ correction in MHz/ Z^6



Tests

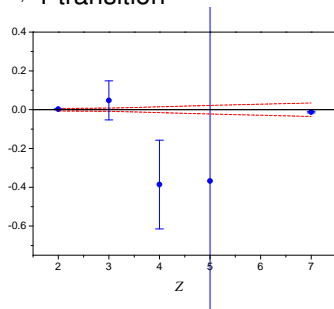
- checking the hydrogenic limit



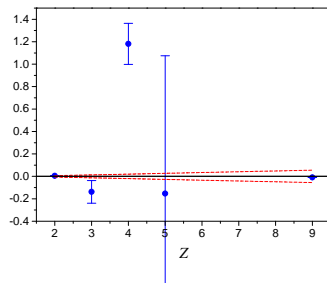
- comparison with experiment for different nuclear charges

Differences: theory-experiment in kHz/Z^8

$0 \rightarrow 1$ transition



$1 \rightarrow 2$ transition



$Z = 3$ Riis *et al.* 1994

$Z = 4$ Scholl *et al.* 1993

$Z = 5$ Dinneen *et al.* 1991

$Z = 7$ Thompson *et al.* 1998

$Z = 9$ Myers *et al.* 1999

Conclusions

- $\alpha^{-1}(\text{He}) = 137.035\,999\,5(39)_{\text{theo}}(6)_{\text{exp}}$ from Shiner 2010
- disagreement between experimental results for ν_{12} and ν_{02}
- large uncertainty due to higher order terms
- testing $1/Z$ expansion against hydrogenic limit
- possible 10^{-9} determination of α requires more accurate estimation of higher order terms
- one of the most accurate tests of QED, ($\mu\text{H}+\text{H}$ Lamb shift)